<u>Title</u>: Global Existence and Asymptotic Behavior for the Compressible Navier-Stokes Equations with a Non-autonomous External Force and a Heat Source

Authors:

Yuming QIN
Department of Applied Mathematics
College of Science
Donghua University
Shanghai 200051, P. R. China
E-mail: yuming_qin@hotmail.com

and

The Institute of Mathematical Sciences The Chinese University of Hong Kong Shatin, N. T., Hong Kong E-mail: ymqin@ims.cuhk.edu.hk

and

Xiaona Yu
Department of Applied Mathematics
College of Science
Donghua University
Shanghai 200051, P. R. China
E-mail: balinda@mail.dhu.edu.cn

Proposed running head:

Global Existence and Asymptotic Behavior for the Compressible Navier-Stokes Equations

Mailing address:

Yuming Qin
Department of Applied Mathematics
College of Science
Donghua University, Shanghai 200051
P. R. China
E-mail: yuming_qin@hotmail.com

Global existence and asymptotic behavior for the compressible Navier-Stokes equations with a non-autonomous external force and a heat source¹

 $Yuming\ Qin^{1,2}\ Xiaona\ Yu^1$

- 1. Department of Applied Mathematics, College of Science, Donghua University, Shanghai 200051, P.R. China. E-mails: yuming_qin@hotmail.com;balinda@mail.dhu.edu.cn
- 2. The Institute of Mathematical Sciences, The Chinese University of Hong Kong, Shatin, N. T., Hong Kong. E-mail: ymqin@ims.cuhk.edu.hk

Abstract

In this paper, we prove the global existence and asymptotic behavior, as time tends to infinity, of solutions in H^i (i=1,2) to the initial boundary value problem of the compressible Navier-Stokes equations of one-dimensional motion of a viscous heat conducting gas in a bounded region with a non-autonomous external force and a heat source. Some new ideas and more delicate estimates are used to prove these results.

Key words: global existence, asymptotic behavior, a uniform priori estimates

1 Introduction

This paper is concerned with global existence and asymptotic behavior to the equations of one-dimensional motion of a viscous heat conducting gas in a bounded region with a non-autonomous external force and a heat source. The Lagrangian form of the conservation laws of mass, momentum, and energy for a one-dimensional gas is (see, e.g., [1, 29, 30])

$$u_t - v_x = 0, (1.1)$$

$$v_t = \sigma_x + f(\int_0^x u dy, t), \tag{1.1}$$

$$v_t = \sigma_x + f(\int_0^x u dy, t), \tag{1.2}$$

$$e_t = -Q_x + \sigma v_x + g(\int_0^x u dy, t). \tag{1.3}$$

Here subscripts indicate partial differentiations, u, v, σ, e, Q and θ denote specific volume, velocity, stress, internal energy, heat flux and absolute temperature, respectively. f, g denote the non-autonomous external force and the heat source, respectively.

For simplicity, we only consider the polytropic viscous ideal gas, i.e.,

$$e = c_v \theta, \quad \sigma = -p + \mu \frac{v_x}{u}, \quad Q = -k \frac{\theta_x}{u}, \quad p = R \frac{\theta}{u}$$
 (1.4)

¹This work will appear in Math. Meth. Appl. Sci.

with suitable positive constants c_v , R, μ and k.

We consider problem (1.1)–(1.3) in the region $\{0 \le x \le 1, t \ge 0\}$ under initial conditions

$$u(x,0) = u_0(x), v(x,0) = v_0(x), \theta(x,0) = \theta_0(x) \text{ on } [0,1]$$
 (1.5)

and the boundary conditions of the form

$$v(0,t) = v(1,t) = 0, Q(0,t) = Q(1,t) = 0, \quad \forall t \ge 0.$$
(1.6)

Throughout the paper we suppose:

$$\int_0^1 u_0(x)dx \equiv \overline{u_0},\tag{1.7}$$

$$0 < C_0^{-1} \le u_0(x) \le C_0, \quad \forall \quad 0 \le x \le 1$$
 (1.8)

where C_0 is a positive constant. Furthermore, we assume that for any $u(x, \cdot) \in L^{\infty}(R^+, L^1[0, 1])$ with $\xi(x,t) \equiv \int_0^x u(y,t)dy$ and $\hat{f}(x,t) \equiv \int_0^t f(\int_0^x u(y,s)dy,s)ds) = \int_0^t f(\xi(x,s),s)ds$, the non-autonomous external force $f = f(\xi(x,t),t)$ and heat source $g = g(\xi(x,t),t)$ satisfy the following conditions

$$f(\xi(x,\cdot),\cdot) \in L^{\infty}(R^{+}, L^{2}[0,1]) \cap L^{2}(R^{+}, L^{\infty}[0,1]) \cap L^{1}(R^{+}, L^{1}[0,1]),$$

$$\hat{f}(x,\cdot) \in L^{1}(R^{+}, L^{2}[0,1]), f_{\xi}(\xi(x,\cdot),\cdot) \in L^{2}(R^{+}, L^{2}[0,1]),$$

$$f_{t}(\xi(x,\cdot),\cdot) \in L^{2}(R^{+}, L^{2}[0,1]), \qquad (1.9)$$

$$g(\xi,t) > 0, g(\xi(x,\cdot),\cdot) \in L^{\infty}(R^{+}, L^{2}[0,1]) \cap L^{2}(R^{+}, L^{2}[0,1]) \cap L^{1}(R^{+}, L^{\infty}[0,1]),$$

$$g_{\xi}(\xi(x,\cdot),\cdot) \in L^{2}(R^{+}, L^{2}[0,1]). \qquad (1.10)$$

Before stating and proving our results, let us first recall the related results in the literature.

For the case of ideal gas (1.4) or real gas with $f \equiv 0, g \equiv 0$, the global existence and asymptotic of smooth (generalized) solutions to the system have been investigated by many authors, e.g., see [2-7, 9-15, 24] on the initial boundary value problems and the Cauchy problem. Moreover, Zheng and Qin [28] obtained the existence of maximal attractor for the problem (1.1)-(1.3) and (1.5),(1.6) for $f \equiv 0, g \equiv 0$. Qin et al [24] established the global existence and large-time behavior of solutions in $H^i(i=2,4)$ for the Cauchy problem. Qin [19] established the existence and exponential stability of a C_0 -semigroup in the subspace of $H^i \times H^i \times H^i$ (i=1,2) for a viscous ideal gas in a bounded domain in R and in a bounded annular domain $G_n = \{x \in R^n | 0 < a < |x| < b\}$ (n=2,3) in R^n for a viscous spherically symmetric ideal gas. This result improved those in [13] for an ideal gas and in [3] for the viscous spherically symmetric ideal gas in G_n . As it is known, the constitutive equations of a real gas are well approximated within moderate ranges of u and θ by the model of an ideal gas. However, under very high temperatures, it becomes inadequate.

For $f \neq 0$, $g \equiv 0$, when the temperature θ is a constant, Qin and Zhao [23] have proved the global existence and asymptotic behavior of solutions in H^2 , Piotr Boguslaw Mucha [8]

obtained the exponential stability under various boundary conditions. Shigemori Yanagi [26] have proved the existence of classical solutions. Zhang and Fang [27] also obtained the global existence, asymptotic behavior, exponential stability with the free boundary condition for the isentropic compressible Navier-Stokes equations. Moreover, we would like to refer the works in [16-30] for the related models.

The methods used in this paper come from those in Qin [16-24], in which the global existence, asymptotic behavior, exponential stability and the existence of universal attractor of solutions to equations for a nonlinear one-dimensional viscous heat-conducting real gas were proved.

The notation in this paper will be as follows:

 $L^{\bar{p}}, 1 \leq \bar{p} \leq +\infty, W^{m,\bar{p}}, m \in N, H^1 = W^{1,2}, H^1_0 = W^{1,2}_0$ denote the usual (Sobolev) spaces on (0,1). In addition, $\|\cdot\|_B$ denotes the norm in the space B; we also put $\|\cdot\| = \|\cdot\|_{L^2}$. We denote by $C^k(I,B), k \in N_0$, the space of k-times continuously differentiable functions from $I \subseteq R$ into a Banach space B, and likewise by $L^{\bar{p}}(I,B), 1 \leq \bar{p} \leq +\infty$ the corresponding Lebesgue spaces. Subscripts t and x denote the (partial) derivatives with respect to t and x, respectively. We use $C_i(i=1,2)$ to denote the generic positive constant depending only on the H^i -norm (i=1,2) of the initial data and $\min_{x \in [0,1]} \theta_0(x), \min_{x \in [0,1]} u_0(x)$.

We now in a position to state our main theorems.

Theorem 1.1 Assume conditions (1.5)-(1.10) hold. Then for any $(u_0, v_0, \theta_0) \in H^1[0, 1] \times H^1_0[0, 1] \times H^1[0, 1]$ with $\theta_0(x) > 0$ for any $x \in [0, 1]$, and that the compatibility conditions hold. Then problem (1.1)-(1.4) admits a unique, uniformly bounded, global-in-time solution $(u, v, \theta) \in H^1[0, 1] \times H^1_0[0, 1] \times H^1[0, 1]$ such that

$$0 < C_1^{-1} \le u(x,t) \le C_1, \forall (x,t) \in [0,1] \times [0,+\infty),$$

$$\|u(t)\|_{H^1}^2 + \|v(t)\|_{H^1}^2 + \|\theta(t)\|_{H^1}^2 + \int_0^t (\|u_x\|^2 + \|v_x\|^2 + \|v_t\|^2 + \|v_t\|^2 + \|v_x\|^2 + \|\theta_x\|^2 + \|\theta_t\|^2 + \|\theta_t\|^2 + \|\theta_{xx}\|^2)(s)ds \le C_1, \quad \forall t > 0.$$

$$(1.11)$$

Moreover, as $t \to +\infty$, we have

$$||u(t) - \bar{u}||_{H^1} \longrightarrow 0, ||v(t)||_{H^1} \longrightarrow 0,$$

$$||\theta(t) - \bar{\theta}||_{H^1} \longrightarrow 0, ||\theta(t) - \bar{\theta}||_{L^{\infty}} \longrightarrow 0$$
(1.13)

where $\bar{u} = \int_0^1 u(x,t) dx = \int_0^1 u_0(x) dx, \bar{\theta} = \int_0^1 \theta(x,t) dx.$

Theorem 1.2 Assume conditions (1.5)-(1.10) hold. Then for any $(u_0, v_0, \theta_0) \in H^2[0, 1] \times H^2_0[0, 1] \times H^2[0, 1]$, problem (1.1)-(1.4) admits a unique, uniformly bounded, global-in-time solution $(u, v, \theta) \in H^2[0, 1] \times H^2_0[0, 1] \times H^2[0, 1]$ such that

$$||u(t)||_{H^{2}}^{2} + ||v(t)||_{H^{2}}^{2} + ||\theta(t)||_{H^{2}}^{2} + \int_{0}^{t} (||u_{x}||_{H^{1}}^{2} + ||v_{x}||_{H^{2}}^{2} + ||v_{t}||_{H^{1}}^{2} + ||\theta_{t}||_{H^{1}}^{2})(s)ds \le C_{1}, \quad \forall t > 0.$$

$$(1.14)$$

Moreover, as $t \to +\infty$, we have

$$||u(t) - \bar{u}||_{H^2} \longrightarrow 0, ||v(t)||_{H^2} \longrightarrow 0, ||\theta(t) - \bar{\theta}||_{H^2} \longrightarrow 0$$
 (1.15)

where $\bar{u} = \int_0^1 u(x,t) dx = \int_0^1 u_0(x) dx, \bar{\theta} = \int_0^1 \theta(x,t) dx.$

We organize our present paper as follows. We will complete the proofs of Theorem 1.1 in Section 2 and Section 3. In Section 4 and Section 5 we will complete the proof of Theorem 1.2.

2 Global Existence in H^1

The proof of global existence of a solution in Theorem 1.1 is based on a priori estimates that can be used to continue a local solution globally in time. The existence and uniqueness of local solutions (with positive u and θ) can be obtained by the linearization of the problem (1.1)-(1.6), and by use of the Banach contraction mapping theorem (cf. [24]). In this section, we assume that assumptions in Theorem 1.1 hold.

We begin with the following lemma.

Lemma 2.1 The following estimates hold,

$$\theta(x,t) > 0, \ \forall (x,t) \in [0,1] \times [0,\infty),$$
 (2.1)

$$\int_{0}^{1} u(x,t)dx = \int_{0}^{1} u_{0}(x)dx \equiv \overline{u_{0}}, \quad \forall t > 0,$$
(2.2)

$$\int_0^1 \left(v^2 + c_v(\theta - \log \theta - 1) + R(u - \log u - 1) \right) dx + C_v(\theta - \log \theta - 1) + R(u - \log u - 1) dx + C_v(\theta - \log \theta - 1) + R(u - \log u - 1) dx + C_v(\theta - \log \theta - 1) + R(u - \log u - 1) dx + C_v(\theta - \log \theta - 1) dx +$$

$$\int_{0}^{t} \int_{0}^{1} \left(\frac{v_x^2}{u\theta} + \frac{\theta_x^2}{u\theta^2} + \frac{g}{\theta} \right) dx ds \le e_0, \quad \forall t > 0.$$
 (2.3)

Proof Inequality (2.1) is a consequence of the generalized maximum principle applied to the following equation, which is equivalent to (1.3),

$$c_v \theta_t - k \left[\frac{\theta_x}{u}\right]_x + \frac{R\theta v_x}{u} = \frac{\mu v_x^2}{u} + g\left(\int_0^x u dy, t\right)$$
 (2.4)

by considering the positivity of θ_0 and g in (1.10).

Integrating (1.1) with respect to x and t, by conditions (1.6)-(1.7), we can easily get (2.2).

Multiplying (1.2) by v, combining with (1.3) and integrating the resultant over [0,1], we can get

$$\int_{0}^{1} (vv_{t} + c_{v}\theta_{t})dx = \int_{0}^{1} (\sigma_{x}v + \sigma v_{x} + k(\frac{\theta_{x}}{u})_{x} + fv + g)dx.$$
 (2.5)

That is,

$$\frac{d}{dt} \int_0^1 (c_v \theta + \frac{v^2}{2}) dx = \int_0^1 (k(\frac{\theta_x}{u})_x + (\sigma v)_x + fv + g) dx = \int_0^1 (fv + g) dx.$$
 (2.6)

Multiplying (1.3) by θ^{-1} and integrating the resultant over [0, 1], we can get

$$\int_0^1 c_v \frac{\theta_t}{\theta} dx = \int_0^1 \left[k \left(\frac{\theta_x}{u} \right)_x \frac{1}{\theta} + \frac{\sigma v_x}{\theta} + \frac{g}{\theta} \right] dx. \tag{2.7}$$

That is,

$$\frac{d}{dt} \int_0^1 (c_v \log \theta + R \log u) dx = \int_0^1 \left(k \frac{\theta_x^2}{u\theta^2} + \mu \frac{v_x^2}{u\theta} + \frac{g}{\theta}\right) dx. \tag{2.8}$$

Adding (2.6) and (2.8) up, we get

$$\frac{d}{dt} \int_{0}^{1} (c_{v}\theta + \frac{v^{2}}{2} - c_{v} \log \theta - R \log u) dx + \int_{0}^{1} (k \frac{\theta_{x}^{2}}{u\theta^{2}} + \mu \frac{v_{x}^{2}}{u\theta} + \frac{g}{\theta}) dx
= \int_{0}^{1} (fv + g) dx.$$
(2.9)

Integrating (2.9) over $[0,t] \times [0,1]$, and using (2.2) and estimates

$$\theta - 1 - \log \theta \ge 0, \quad for \quad \theta > 0, \tag{2.10}$$

$$u - 1 - \log u \ge 0, \text{ for } u > 0,$$
 (2.11)

we have

$$\int_{0}^{1} \left[\frac{1}{2} v^{2} + c_{v} (\theta - \log \theta - 1) + R(u - \log u - 1) \right] dx
+ \int_{0}^{t} \int_{0}^{1} \left[\frac{k \theta_{x}^{2}}{u \theta^{2}} + \frac{\mu v_{x}^{2}}{u \theta} + \frac{g}{\theta} \right] dx ds
\leq C_{1} + \int_{0}^{t} \int_{0}^{1} (f v + g) dx ds
\leq C_{1} + \frac{1}{2} \int_{0}^{t} \int_{0}^{1} ||f||_{L^{2}[0,1]} (||v||^{2} + 1) ds + \int_{0}^{t} ||g||_{L^{1}[0,1]} ds
\leq C_{1} + \frac{1}{2} \int_{0}^{t} ||f||_{L^{2}[0,1]} ||v||^{2} ds.$$
(2.12)

Applying the Gronwall inequality to (2.12) and using (1.9) yields the desired estimate (2.3). $\hfill\Box$

The next lemma was given when $f = g \equiv 0$ in [2, 14-17] for a viscous heat-conducting real gas and in [6] for a viscous polytropic ideal gas.

Lemma 2.2 For any $t \ge 0$, there exists one point $x_1 = x_1(t) \in [0,1]$ such that the solution u(x,t) to problem (1.1)-(1.6) possesses the following expression:

$$u(x,t) = D(x,t)Z(t)\left[1 + \frac{R}{\mu} \int_0^t \frac{\theta(x,s)}{D(x,s)Z(s)} ds\right]$$
(2.13)

where

$$D(x,t) = u_0(x) \exp\left\{\frac{1}{\mu} \left[\int_{x_1(t)}^x v dy - \int_0^x v_0 dy + \frac{1}{\overline{u_0}} \int_0^1 (u_0 \int_0^x v_0 dy) dx \right] + \frac{1}{\overline{u_0}} \int_0^t \int_0^1 v(x,s) \int_0^s f(\int_0^x u(z,\tau) dz,\tau) d\tau dx ds - \int_{x_1(t)}^x \int_0^t f(\int_0^y u(z,s) dz,s) ds dy \right\},$$
(2.14)

$$Z(t) = \exp\left[-\frac{1}{\mu\overline{u_0}} \int_0^t \int_0^1 (v^2 + R\theta) dx ds\right]. \tag{2.15}$$

Proof Let

$$h(x,t) = \int_0^x v_0(y)dy + \int_0^t \sigma(x,s)ds.$$
 (2.16)

Then we infer from (1.1)-(1.2)

$$h_x = v_0(x) + \int_0^t \sigma_x(x,s)ds = v(x,t) - \int_0^t f(\int_0^x u(y,\tau)dy,\tau)ds,$$
 (2.17)

$$h_t = \sigma(x,t) = -p + \mu \frac{v_x}{u}.$$
 (2.18)

By (1.1) and (2.18), we get

$$(uh)_t = u_t h + uh_t = v_x h - up + \mu v_x. \tag{2.19}$$

Integrating (2.19) over $[0,1] \times [0,t]$, by (1.6), we can conclude

$$\int_{0}^{1} uhdx = \int_{0}^{1} u_{0}h_{0}dx + \int_{0}^{t} \int_{0}^{1} (v_{x}h - up + \mu v_{x})dxds
= \int_{0}^{1} u_{0}h_{0}dx - \int_{0}^{t} \int_{0}^{1} [v^{2} + up - v \int_{0}^{s} fd\tau](x, s)dxds.$$
(2.20)

Applying the mean value theorem to (2.20) and using (2.2), we derive that there exists one point $x_1 = x_1(t) \in [0, 1]$ such that

$$\int_{0}^{1} uhdx = h(x_{1}(t), t) \int_{0}^{1} u(x, t)dx = \overline{u_{0}}h(x_{1}(t), t)$$
(2.21)

with $\overline{u_0} = \int_0^1 u_0(x) dx$. Thus it follows from (2.16) and (2.21) that

$$\int_{0}^{t} \sigma(x_{1}(t), s) ds = h(x_{1}(t), t) - \int_{0}^{x_{1}(t)} v_{0}(y) dy$$

$$= \frac{1}{\overline{u_{0}}} \int_{0}^{1} h(x, t) u(x, t) dx - \int_{0}^{x_{1}(t)} v_{0}(y) dy$$

$$= \frac{1}{\overline{u_{0}}} \left[\int_{0}^{1} u_{0}(x) h_{0}(x) dx - \int_{0}^{t} \int_{0}^{1} (v^{2} + up)(x, s) dx ds \right]$$

$$+ \int_{0}^{t} \int_{0}^{1} v(x, s) \int_{0}^{s} f(\int_{0}^{x} u(z, \tau) dz, \tau) d\tau dx ds - \int_{0}^{x_{1}(t)} v_{0}(y) dy. \tag{2.22}$$

By (1.1) and (1.4), we can rewrite (1.2) as

$$v_t - \mu[\log u]_{xt} = -(\frac{R\theta}{u})_x + f.$$
 (2.23)

Integrating (2.23) over $[x_1(t), x] \times [0, t]$, we can infer

$$\mu \log u - R \int_0^t \frac{\theta}{u} ds = \mu \log u_0 + \int_0^t \sigma(x_1(t), s) ds + \int_{x_1(t)}^x (v - v_0) dy - \int_{x_1(t)}^x \int_0^t f ds dy. \quad (2.24)$$

Inserting (2.22) in (2.24), we obtain

$$\mu \log u - R \int_{0}^{t} \frac{\theta}{u} ds = \mu \log u_{0} - \frac{1}{\overline{u_{0}}} \int_{0}^{t} \int_{0}^{1} (v^{2} + up) dx ds + \frac{1}{\overline{u_{0}}} \int_{0}^{1} u_{0} h_{0} dx + \frac{1}{\overline{u_{0}}} \int_{0}^{t} \int_{0}^{1} v \int_{0}^{s} f(\int_{0}^{x} u(z, \tau) dz, \tau) d\tau ds dx + \int_{x_{1}(t)}^{x} v(y, t) dy - \int_{0}^{x} v_{0}(y) dy - \int_{x_{1}(t)}^{x} \int_{0}^{t} f ds dy.$$

$$(2.25)$$

With the definitions of D(x,t) and Z(t), we infer from (2.25) that

$$\frac{Z(t)}{B(x,t)} = \frac{1}{u(x,t)} \exp\left[\frac{R}{\mu} \int_0^t \frac{\theta(x,s)}{u(x,s)} ds\right]. \tag{2.26}$$

Multiplying (2.26) by $R\theta/u$, and integrating the result with respect to t, we have

$$\exp\left[\frac{R}{\mu} \int_0^t \frac{\theta(x,s)}{u(x,s)} ds\right] = 1 + \frac{R}{\mu} \int_0^t \frac{\theta(x,s)Z(s)}{D(x,s)} ds$$

which along with (2.26) gives (2.13).

Lemma 2.3

$$0 < C_1^{-1} \le u(x,t) \le C_1, \ \forall (x,t) \in [0,1] \times [0,+\infty).$$
 (2.27)

Proof Let

$$M_u(t) = \max_{x \in [0,1]} u(x,t), \quad m_u(t) = \min_{x \in [0,1]} u(x,t),$$

 $M_{\theta}(t) = \max_{x \in [0,1]} \theta(x,t), \quad m_{\theta}(t) = \min_{x \in [0,1]} \theta(x,t).$

It follows from (2.3) and the convexity of the function $-\log y$ that

$$\int_{0}^{1} \theta dx - \log \int_{0}^{1} \theta dx - 1 \le \int_{0}^{1} (\theta - \log \theta - 1) dx \le C_{1}$$

which implies that there exist a point $b(t) \in [0,1]$ and two positive constants r_1, r_2 such that

$$0 < r_1 \le \int_0^1 \theta(x, t) dx = \theta(b(t), t) \le r_2$$
 (2.28)

with r_1, r_2 being two positive roots of the equation $y - \log y - 1 = C_1$.

By Young's inequality, Hölder's inequality, (1.7), (1.9) and Lemma 2.1, we have

$$\left| \int_{x_1}^x v dy - \int_0^x v_0 dy + \frac{1}{\overline{u_0}} \int_0^1 (u_0 \int_0^x v_0 dy) dx \right|$$

$$\leq \|v\|_{L^1[0,1]} + 2\|v_0\|_{L^1[0,1]} \leq C_1,$$
(2.29)

$$\left| \int_{x_1(t)}^x \int_0^t f(\int_0^y u(z,s)dz,s)dsdy \right| \le ||f||_{L^1(R^+,L^1[0,1])} \le C_1, \tag{2.30}$$

$$\left| \frac{1}{\overline{u_0}} \int_0^t \int_0^1 v(x,s) \int_0^s f(\int_0^x u(z,\tau) dz,\tau) d\tau dx ds \right| \le C_1 \int_0^t \|v(x,s)\| \|\hat{f}(x,s)\| ds$$

$$\le C_1 \int_0^t \|\hat{f}(x,s)\| ds \le C_1.$$
(2.31)

Thus we easily deduce from Lemmas 2.1-2.2, (2.28)-(2.30) that

$$0 < C_1^{-1} \le D(x,t) \le C_1, \ \forall (x,t) \in [0,1] \times [0,+\infty).$$
 (2.32)

On the other hand, we readily infer from (2.3), (2.28) that for $0 \le s \le t$,

$$Rr_1(t-s) \le R \int_s^t \int_0^1 \theta dx d\tau \le \int_s^t \int_0^1 (v^2 + R\theta)(x, \tau) dx d\tau$$

 $\le \int_s^t [\|v\|^2 + \int_0^1 R\theta dx](\tau) d\tau$
 $\le (e_0 + Rr_2)(t-s)$

which, combined with (2.15), implies that for any $0 \le s \le t$,

$$C_1^{-1}e^{-a_2(t-s)} \le Z(t)Z^{-1}(s) = \exp\left[-\frac{1}{\mu\overline{u_0}} \int_s^t \int_0^1 (v^2 + R\theta) dx ds\right]$$

$$\le C_1 e^{-a_1(t-s)}$$
(2.33)

with $a_1 = \frac{r_1 R}{\mu \overline{u_0}}$ and $a_2 = \frac{r_2 R + e_0}{\mu \overline{u_0}}$. We have from (2.28)

$$\begin{aligned} |\theta^{1/2}(x,t) - \theta^{1/2}(b(t),t)| & \leq C \left| \int_{b(t)}^{x} \theta^{-1/2} \theta_x dx \right| \leq C \left(\int_{0}^{1} \frac{\theta_x^2}{u\theta} dx \right)^{\frac{1}{2}} \left(\int_{0}^{1} u \theta^2 dx \right)^{\frac{1}{2}} \\ & \leq C V^{\frac{1}{2}}(t) M_u^{\frac{1}{2}}(t) \end{aligned}$$

where $V(t) = \int_0^1 \frac{\theta_x^2}{u\theta^2} dx$ satisfies $\int_0^\infty V(s) ds < +\infty$, by (2.3) in Lemma 2.1, which gives

$$C_1^{-1} - C_1 V(t) M_u(t) \le \theta(x, t) \le C_1 + C V(t) M_u(t).$$
(2.34)

Thus we conclude from Lemma 2.2 and (2.28)-(2.34)

$$u(x,t) = D(x,t) \left[Z(t) + \frac{R}{\mu} \int_0^t (\theta(x,s)D^{-1}(x,s)Z(t)Z^{-1}(s)ds \right]$$

$$\leq C_1 \left[e^{-a_1t} + \int_0^t (1+V(s)M_u(s))e^{-a_1(t-s)}ds \right]$$

$$\leq C_1 + C_1 \int_0^t M_u(s)V(s)ds$$

whence

$$M_u(t) \le C_1 + C_1 \int_0^t M_u(s)V(s)ds$$

which, by using Gronwall's inequality and (2.3), implies

$$M_u(t) \le C_1. \tag{2.35}$$

So, from (2.34)-(2.35),

$$C_1^{-1} - C_1 V(t) \le \theta(x, t) \le C_1 + C_1 V_1(t).$$
 (2.36)

Similarly, noting that as $t \to +\infty$,

$$\int_0^t V(s)e^{-a_2(t-s)}ds \le e^{-a_2t/2} \int_0^{t/2} V(s)ds + \int_{t/2}^t V(s)ds \to 0,$$

we infer from (2.31)-(2.36) and Lemma 2.2 that there exists a large time $t_0 > 0$ such that as $t \ge t_0$, $\forall x \in [0, 1]$,

$$u(x,t) = D(x,t) \left[Z(t) + \frac{R}{\mu} \int_0^t \theta D^{-1}(x,s) Z(t) Z^{-1}(s) ds \right]$$

$$\geq C_1^{-1} \left[e^{-a_2 t} + \int_0^t (C_1^{-1} - C_1 V(s)) e^{-a_2 (t-s)} ds \right]$$

$$\geq C_1^{-1} - C_1 \int_0^t V(s) e^{-a_2 (t-s)} ds \geq (2C_1)^{-1}. \tag{2.37}$$

Noting that $D(x,t) \ge 1$, $Z(t) \ge 1$ and $\theta(x,t) > 0$, we infer from Lemma 2.2 that

$$u(x,t) \ge D(x,t)Z(t) \ge C_1^{-1}e^{-a_2t} \ge C_1^{-1}e^{-a_2t_0}, \ \forall (x,t) \in [0,1] \times [0,t_0]$$

which with (2.37) and (2.35) gives (2.27).

Corollary 2.1 There holds that for any t > 0,

$$\int_0^t \|v(s)\|_{L^{\infty}[0,1]}^2 \le C_1. \tag{2.38}$$

Proof By Lemmas 2.1-2.4, we can derive

$$\int_{0}^{t} \|v(s)\|_{L^{\infty}[0,1]}^{2} ds \leq \int_{0}^{t} (\int_{0}^{1} |v_{x}| dx)^{2} ds
\leq \int_{0}^{t} \int_{0}^{1} \frac{v_{x}^{2}}{\theta} dx \int_{0}^{1} \theta dx ds \leq C_{1}.$$

Lemma 2.4 The following estimate holds

$$\|\theta(t)\|^{2} + \|v(t)\|_{L^{4}[0,1]}^{4} + \int_{0}^{t} (\|\theta_{x}\|^{2} + \|vv_{x}\|^{2})(s)ds \le C_{1}, \quad \forall t > 0.$$
 (2.39)

Proof Multiplying (2.5) by $c_v\theta + \frac{v^2}{2}$, and integrating the resulting equation by parts, we arrive at

$$\frac{1}{2}\frac{d}{dt}\|c_v\theta + \frac{v^2}{2}\|^2 + \int_0^1 (kc_v\frac{\theta_x^2}{u} + \mu\frac{v^2v_x^2}{u})dx
= \int_0^1 \left[c_vpv\theta_x + pv^2v_x - \mu\frac{c_vvv_x\theta_x}{u} - k\frac{\theta_xvv_x}{u} + (vf + g)(c_v\theta + \frac{v^2}{2})\right]dx. \quad (2.40)$$

9

Integrating (2.40) over [0,t], by Lemma 2.3 and Young's inequality, we can obtain

$$||c_{v}\theta + \frac{v^{2}}{2}||^{2} + \int_{0}^{t} (||vv_{x}||^{2} + ||\theta_{x}||^{2})(s)ds$$

$$\leq C_{1} \Big[\int_{0}^{t} \int_{0}^{1} (|pv(\theta_{x}| + |vv_{x}|) + |vv_{x}\theta_{x}| + |v\theta f| + |v^{3}f| + |g\theta| + |v^{2}g|)dxds \Big]$$

$$\leq \varepsilon \int_{0}^{t} ||\theta_{x}(s)||^{2}ds + C_{1} \int_{0}^{t} [||vv_{x}||^{2} + ||\theta v||^{2} + ||f||^{2} + ||g||_{L^{\infty}[0,1]} (||\theta||_{L^{1}[0,1]} ||v||^{2})](s)ds$$

$$+ \int_{0}^{t} (||f||_{L^{\infty}[0,1]}^{2} ||v||_{L^{4}[0,1]}^{4} + ||v||^{2})(s)ds.$$

$$(2.41)$$

Thus, by (1.7)-(1.9) and Corollary 2.1, for $\varepsilon > 0$ small enough,

$$\|\theta(t)\|^{2} + \|v(t)\|_{L^{4}[0,1]}^{4} + \int_{0}^{t} \|\theta_{x}(s)\|^{2} ds$$

$$\leq C_{1} + C_{1} \int_{0}^{t} [\|vv_{x}\|^{2} + \|\theta\|^{2} \|v\|_{L^{\infty}[0,1]}^{2} + \|f\|_{L^{\infty}[0,1]}^{2} \|v\|_{L^{4}[0,1]}^{4}](s) ds. \quad (2.42)$$

Multiplying (1.2) by v^3 , and integrating it over [0, 1], we can conclude

$$\frac{d}{dt} \int_0^1 \frac{v^4}{4} dx = \int_0^1 (3v^2 v_x (R\frac{\theta}{u} - \mu \frac{v_x}{u}) + fv^3) dx.$$
 (2.43)

Integrating (2.43) over [0,t] and by Young's inequality, we can get

$$||v(t)||_{L^{4}[0,1]}^{4} + \int_{0}^{t} ||vv_{x}||^{2}(s)ds \leq C_{1} + C_{1} \int_{0}^{t} \int_{0}^{1} (|\theta v^{2}v_{x}| + |fv^{3}|) dxds$$

$$\leq C_{1} + \frac{1}{2} \int_{0}^{t} ||vv_{x}||^{2}(s) ds + C_{1} \int_{0}^{t} [||v||_{L^{\infty}[0,1]}^{2} ||\theta||^{2} + ||f||_{L^{\infty}[0,1]}^{2} ||v||_{L^{4}[0,1]}^{4} + ||v||^{2}](s) ds$$

whence

$$||v(t)||_{L^{4}[0,1]}^{4} + \int_{0}^{t} ||vv_{x}||^{2}(s)ds \le C_{1} + C_{1} \int_{0}^{t} [||v||_{L^{\infty}[0,1]}^{2} ||\theta||^{2} + ||f||_{L^{\infty}[0,1]}^{2} ||v||_{L^{4}[0,1]}^{4}](s)ds.$$

$$(2.44)$$

Multiplying (2.44) by a large constant and adding it to (2.42) yields

$$\|\theta(t)\|^{2} + \|v(t)\|_{L^{4}[0,1]}^{4} + \int_{0}^{t} (\|\theta_{x}\|^{2} + \|vv_{x}\|^{2})(s)ds$$

$$\leq C_{1} + C_{1} \int_{0}^{t} [\|v\|_{L^{\infty}[0,1]}^{2} + \|f\|_{L^{\infty}[0,1]}^{2}] (\|\theta\|^{2} + \|v\|_{L^{4}[0,1]}^{4})ds. \tag{2.45}$$

Applying the Gronwall inequality to (2.45), and using Corollary 2.1, (1.7) and the following estimates,

$$\int_{0}^{t} \|f\|_{L^{\infty}[0,1]}^{2} ds \leq C \int_{0}^{t} (\|f\| \|f_{x}\| + \|f\|^{2}) ds \leq C \int_{0}^{t} (\|f\|^{2} + \|f_{x}\|^{2})
\leq C \int_{0}^{t} (\|f\|^{2} + \|f_{\xi}\|^{2}) \leq C_{1},
\int_{0}^{t} \|g\|_{L^{\infty}[0,1]} ds \leq C \int_{0}^{t} (\|g\|^{1/2} \|g_{x}\|^{1/2} + \|g\|) ds \leq C \int_{0}^{t} (\|g\| + \|g_{x}\|)
\leq C \int_{0}^{t} (\|g\| + \|g_{\xi}\|) \leq C_{1},$$

we can conclude (2.39).

Lemma 2.5 The following estimates holds

$$||u_x(t)||^2 + \int_0^t \int_0^1 \theta u_x^2 dx ds \le C_1, \quad \forall t > 0.$$
 (2.46)

Proof By (1.2), we get

$$\left(v - \mu \frac{u_x}{u}\right)_t - R \frac{\theta u_x}{u^2} = -R \frac{\theta_x}{u} + f. \tag{2.47}$$

Multiplying (2.47) by $v - \mu \frac{u_x}{u}$ and integrating over $[0,1] \times [0,t]$, by Lemmas 2.1-2.4 and Corollary 2.1, we obtain

$$\frac{1}{2}\|v - \mu \frac{u_x}{u}\|^2 + \int_0^t \int_0^1 \mu R \frac{u_x^2 \theta}{u^3} dx ds
= \frac{1}{2}\|v_0 - \mu \frac{u_{0x}}{u}\|^2 + \int_0^t \int_0^1 \left(R \frac{u_x \theta v}{u^2} + (f - R \frac{\theta_x}{u})(v - \mu \frac{u_x}{u})\right) dx ds
\leq C_1 + C_1 \int_0^t \int_0^1 \theta(\varepsilon u_x^2 + C(\varepsilon)v^2) dx ds + C_1 \int_0^t (\|\theta_x\|^2 + \|v\|^2) dx ds
+ \varepsilon \int_0^t \int_0^1 \theta u_x^2 dx ds + C(\varepsilon) \int_0^t \int_0^1 \frac{\theta_x^2}{\theta} dx ds + \int_0^t \int_0^1 f(v - \mu \frac{u_x}{u}) dx ds
\leq C_1 + C_1 \varepsilon \int_0^t \int_0^1 \theta u_x^2 dx ds + C_1 \int_0^t \|v\|_{L^{\infty}[0,1]}^2 \int_0^1 \theta dx ds
+ C_1 \int_0^t \int_0^1 (\frac{\theta_x^2}{\theta^2} + \theta_x^2) dx ds + C_1 \int_0^t (\|f\|^2 + \|v\|^2) ds
+ \varepsilon \sup_{0 \leq s \leq t} \|u_x(s)\| \int_0^t \|f\|_{L^{\infty}[0,1]} ds
\leq C_1 + C_1 \varepsilon \int_0^t \int_0^1 \theta u_x^2 dx ds + \varepsilon \sup_{0 \leq s \leq t} \|u_x(s)\|^2$$

which with (2.2) yields

$$\sup_{0 \le s \le t} \|u_x(s)\|^2 \le C_1 + C_1 \varepsilon \int_0^t \int_0^1 \theta u_x^2 dx ds + \varepsilon \sup_{0 \le s \le t} \|u_x(s)\|^2.$$
 (2.48)

Thus taking $\varepsilon > 0$ sufficiently small in (2.48) gives (2.46).

Lemma 2.6 The following estimates hold,

$$||v(t)||^2 + \int_0^t ||v_x(s)||^2 ds \le C_1, \quad \forall t > 0,$$
(2.49)

$$||v_x(t)||^2 + \int_0^t (||v_{xx}||^2 + ||v_t||^2)(s)ds \le C_1 \left(1 + \sup_{0 \le s \le t} ||\theta(s)||_{L^{\infty}}\right), \quad \forall t > 0. \quad (2.50)$$

Proof Multiplying (1.2) by v, v_{xx} and v_t , respectively, and then integrating the resultants over $[0, 1] \times [0, t]$, using Lemmas 2.1-2.5 and Corollary 2.1, we deduce

$$||v(t)||^{2} + 2\mu \int_{0}^{t} \int_{0}^{1} \frac{v_{x}^{2}}{u} dx ds$$

$$\leq C_{1} + C_{1} \int_{0}^{t} \int_{0}^{1} (-R \frac{\theta_{x} v}{u} + R \frac{\theta u_{x} v}{u^{2}} + f v) dx ds$$

$$\leq C_{1} \int_{0}^{t} (||\theta_{x}||^{2} + ||v||^{2} + \int_{0}^{1} \theta u_{x}^{2} dx + ||v||_{L^{\infty}[0,1]}^{2} \int_{0}^{1} \theta dx + ||f||^{2}) ds$$

$$\leq C_{1}, \tag{2.51}$$

$$\frac{1}{2}\|v_{x}(t)\|^{2} + \mu \int_{0}^{t} \int_{0}^{1} \frac{v_{xx}^{2}}{u} dx ds
= \frac{1}{2}\|v_{0x}\|^{2} + \int_{0}^{t} \int_{0}^{1} \left(R(\frac{\theta}{u})_{x}v_{xx} + \mu \frac{v_{x}u_{x}v_{xx}}{u^{2}} - fv_{xx}\right) dx ds
\leq C_{1} + C_{1} \int_{0}^{t} \int_{0}^{1} \left(|v_{x}u_{x}v_{xx}| + |R(\frac{\theta}{u})_{x}v_{xx}| + |fv_{xx}|\right) dx ds
\leq C_{1} + \varepsilon \int_{0}^{t} \|v_{xx}\|^{2} ds + C_{1} \int_{0}^{t} \left(\|v_{x}\|_{L^{\infty}[0,1]}^{2} \|u_{x}\|^{2} + \|\theta_{x}\|^{2} + \|\theta\|_{L^{\infty}[0,1]} \|\theta^{1/2}u_{x}\|^{2} + \|f\|^{2}\right) ds
\leq \varepsilon \int_{0}^{t} \|v_{xx}\|^{2} ds + C_{1} \int_{0}^{t} \left(\|v_{x}\|\|v_{xx}\|\|u_{x}\|^{2}\right) ds + C_{1} (1 + \sup_{0 \leq s \leq t} \|\theta(s)\|_{L^{\infty}})
\leq 2\varepsilon \int_{0}^{t} \|v_{xx}(s)\|^{2} ds + C_{1} \int_{0}^{t} \|v_{x}\|^{2} ds + C_{1} (1 + \sup_{0 \leq s \leq t} \|\theta(s)\|_{L^{\infty}[0,1]})$$

i.e., for sufficiently small $\varepsilon > 0$,

$$||v_x(t)||^2 + \int_0^t ||v_{xx}(s)||^2 ds \le C_1 (1 + \sup_{0 \le s \le t} ||\theta(s)||_{L^{\infty}[0,1]})$$
(2.52)

and

$$\begin{split} &\|v_x(t)\|^2 + \int_0^t \|v_t(s)\|^2 ds \\ &\leq C_1 + C_1 \int_0^t (\|p_x\|^2 + \|f\|^2 + \|v_x\|_{L^3}^3)(s) ds \\ &\leq C_1 + C_1 \int_0^t (\|\theta u_x\|^2 + \|\theta_x\|^2 + \|v_x\|^{5/2} \|v_{xx}\|^{1/2} + \|f\|^2) ds \\ &\leq C_1 + C_1 \int_0^t [\sup_{0 \leq s \leq t} \|\theta\|_{L^\infty} \int_0^1 \theta u_x^2 dx + \|\theta_x\|^2 + \|f\|^2] ds \\ &+ C_1 \Big(\int_0^t \|v_x\|^{10/3} ds \Big)^{3/4} \Big(\int_0^t \|v_{xx}\|^2 ds \Big)^{1/4} \\ &\leq C_1 \Big(1 + \sup_{0 \leq s \leq t} \|\theta\|_{L^\infty} \Big) + \varepsilon \sup_{0 \leq s \leq t} \|v_x\|^2 \Big(\int_0^t \|v_x\|^2 ds \Big)^{3/2} + C_1(\varepsilon) \Big(\int_0^t \|v_{xx}\|^2 ds \Big)^{1/2} \\ &\leq C_1 \Big(1 + \sup_{0 \leq s \leq t} \|\theta\|_{L^\infty} \Big) + \varepsilon \sup_{0 \leq s \leq t} \|v_x\|^2 \Big(\int_0^t \|v_x\|^2 ds \Big)^{3/2} + C_1(\varepsilon) \Big(\int_0^t \|v_{xx}\|^2 ds \Big)^{1/2} \end{split}$$

which along with (2.52) yields estimate (2.49). Here we have used the following interpolation inequalities,

$$||v_{x}(t)||_{L^{\infty}[0,1]} \leq C||v_{x}(t)||^{1/2}||v_{xx}(t)||^{1/2},$$

$$||v_{x}(t)||_{L^{3}}^{3} \leq C||v_{x}(t)||^{5/2}||v_{xx}(t)||^{1/2} + C||v_{x}(t)||^{3} \leq C||v_{x}(t)||^{5/2}||v_{xx}(t)||^{1/2},$$

$$\bar{v_{x}} = \int_{0}^{1} v_{x} dx = 0.$$

Lemma 2.7 The following estimates hold

$$\|\theta_x(t)\|^2 + \int_0^t \|\theta_{xx}\|^2 ds \le C_1 \left(1 + \sup_{0 \le s \le t} \|\theta(s)\|_{L^{\infty}}\right)^2, \quad \forall t > 0.$$
 (2.53)

Proof Multiplying (1.3) by θ , integrating the resulting equation over $[0,1] \times [0,t]$, we get

$$\frac{1}{2}c_{v}\|\theta_{x}\|^{2} + \int_{0}^{t} \int_{0}^{1} k \frac{\theta_{xx}^{2}}{u} dx ds
= \frac{1}{2}c_{v}\|\theta_{0x}\|^{2} + \int_{0}^{t} \int_{0}^{1} k \frac{\theta_{x}u_{x}\theta_{xx}}{u^{2}} dx ds - \int_{0}^{t} \int_{0}^{1} (\sigma v_{x} + g)\theta_{xx}) dx ds
\leq C_{1} + \varepsilon \int_{0}^{t} \|\theta_{xx}\|^{2} ds + C_{1} \int_{0}^{t} \int_{0}^{1} (\theta_{x}^{2}u_{x}^{2} + \theta^{2}v_{x}^{2} + v_{x}^{4} + g^{2}) dx ds.$$

That is, for small $\varepsilon > 0$,

$$\begin{split} &\|\theta_x(t)\|^2 + \int_0^t \|\theta_{xx}(s)\|^2 ds \\ &\leq C_1 + C_1 \int_0^t \left((\|\theta_x\| \|\theta_{xx}\| + \|\theta_x\|^2) \|u_x\|^2 + \|v_x\|^3 \|v_{xx}\| \right) (s) ds \\ &+ C_1 \sup_{0 \leq s \leq t} \|\theta(s)\|_{L^{\infty}}^2 \int_0^t \|v_x(s)\|^2 ds + \int_0^t \|g\|^2 ds \\ &\leq C_1 (1 + \sup_{0 \leq s \leq t} \|\theta(s)\|_{L^{\infty}})^2 + \frac{1}{2} \int_0^t \|\theta_{xx}(s)\|^2 ds + C_1 \int_0^t \|\theta_x(s)\|^2 (\|u_x\|^4 + \|u_x\|^2) (s) ds \\ &+ C_1 \int_0^t \|v_{xx}(s)\|^2 ds + C_1 \int_0^t \|v_x(s)\|^6 ds \\ &\leq C_1 (1 + \sup_{0 \leq s \leq t} \|\theta(s)\|_{L^{\infty}})^2 + C_1 \int_0^t \|v_{xx}(s)\|^2 ds + \frac{1}{2} \int_0^t \|\theta_{xx}(s)\|^2 ds \\ &+ C_1 \sup_{0 \leq s \leq t} (\|u_x\|^4 + \|u_x\|^2) \int_0^t \|\theta_x(s)\|^2 ds + C_1 \sup_{0 \leq s \leq t} \|v_x(s)\|^4 \int_0^t \|v_x(s)\|^2 ds \\ &\leq C_1 (1 + \sup_{0 \leq s \leq t} \|\theta(s)\|_{L^{\infty}[0,1]})^2. \end{split}$$

Lemma 2.8 The following estimates hold

$$\|\theta(t)\|_{L^{\infty}} \le C_1, \quad \forall t > 0, \tag{2.54}$$

$$\|\theta_{x}(t)\|^{2} + \|v_{x}(t)\|^{2} + \|u_{x}(t)\|^{2} + \int_{0}^{t} (\|\theta_{x}\|^{2} + \|vv_{x}\|^{2} + \|u_{x}\|^{2} + \|v_{x}\|^{2} + \|v_{x}\|^{2} + \|v_{t}\|^{2} + \|v_{t}\|^{2} + \|\theta_{t}\|^{2} + \|\theta_{t}\|^{2})(s)ds \le C_{1}, \quad \forall t > 0.$$
(2.55)

Proof With the help of the Gagliardo-Nirenberg interpolation inequality, Young's inequality and Lemmas 2.1-2.7, we derive

$$\|\theta(t)\|_{L^{\infty}} \leq C\|\theta_{x}(t)\|^{2/3}\|\theta(t)\|_{L^{1}[0,1]}^{1/3} + C\|\theta(t)\|_{L^{1}[0,1]} \leq C(\|\theta_{x}\|^{2/3} + 1)$$

$$\leq C(1 + \sup_{0 \leq s \leq t} \|\theta(s)\|_{L^{\infty}[0,1]}^{2/3})$$

$$\leq \sup_{0 \leq s \leq t} \|\theta(s)\|_{L^{\infty}[0,1]} + C_{1}. \tag{2.56}$$

Taking the supremum on the right-hand side of (2.56), and picking ε sufficiently small, and we can get the inequality (2.54).

Multiplying (1.3) by θ_t , and then integrating the resultants over $[0,1] \times [0,t]$, using Lemmas 2.1-2.7, we infer

$$\begin{split} & \int_{0}^{1} k \frac{\theta_{x}^{2}}{u} dx + \int_{0}^{t} \int_{0}^{1} c_{v} \theta_{t}^{2} dx ds \\ &= \int_{0}^{1} k \frac{\theta_{0x}^{2}}{u_{0}} dx + \int_{0}^{t} \int_{0}^{1} \left(-k \frac{\theta_{x}^{2} v_{x}}{u^{2}} + (-R \frac{\theta}{u} + \mu \frac{v_{x}}{u}) v_{x} \theta_{t} + g \theta_{t} \right) dx ds \\ &\leq C_{1} + C_{1} \int_{0}^{t} \int_{0}^{1} (\theta_{x}^{2} |v_{x}| + \theta |v_{x} \theta_{t}| + v_{x}^{2} |\theta_{t}| + g |\theta_{t}|) dx ds \\ &\leq C_{1} + \varepsilon \int_{0}^{t} \|\theta_{t}(s)\|^{2} ds + C_{1} \int_{0}^{t} (\|v_{x}\|^{2} + \|\theta_{x}\|^{4}_{L^{4}[0,1]} \\ &+ \|\theta\|_{L^{\infty}[0,1]}^{2} \|v_{x}\|^{2} + \|v_{x}\|_{L^{4}[0,1]}^{4} + \|g\|^{2})(s) ds \\ &\leq C_{1} + \varepsilon \int_{0}^{t} \|\theta_{t}(s)\|^{2} ds + C_{1} \int_{0}^{t} (\|\theta_{x}\|^{3} \|\theta_{xx}\| + \|\theta_{x}\|^{4} \\ &+ \|\theta\|_{L^{\infty}[0,1]}^{2} \|v_{x}\|^{2} + \|v_{x}\|^{3} \|v_{xx}\| + \|g\|^{2})(s) ds \\ &\leq C_{1} + \varepsilon \int_{0}^{t} \|\theta_{t}\|^{2} ds + C_{1} \sup_{0 \leq s \leq t} (\|\theta_{x}\|^{2} + \|\theta_{x}\|^{4}) \int_{0}^{t} \|\theta_{x}(s)\|^{2} ds \\ &+ C_{1} \sup_{0 \leq s \leq t} \|v_{x}(s)\|^{4} \int_{0}^{t} \|v_{x}(s)\|^{2} ds + C_{1} \int_{0}^{t} (\|\theta_{xx}\|^{2} + \|v_{x}\|^{2} + \|v_{xx}\|^{2} + \|g\|^{2})(s) ds \\ &\leq C_{1} + \varepsilon \int_{0}^{t} \|\theta_{t}(s)\|^{2} ds \end{split}$$

whence, for small $\varepsilon > 0$,

$$\|\theta_x(t)\|^2 + \int_0^t \|\theta_t(s)\|^2 ds \le C_1$$

which, combined with Lemmas 2.1-2.7 and (2.54), gives the estimate (2.56).

3 Asymptotic Behavior in H^1

In this section we will complete the proof of (1.13) and assume that the assumptions in Theorem 1.1 are valid. We begin with the following lemma.

Lemma 3.1 Suppose that y and h are nonnegative functions on $[0, +\infty)$, y' is locally integrable, and y and h satisfy

$$\forall t \ge 0, \ y'(t) \le A_1 y^2(t) + A_2 + h(t),$$

 $\forall T > 0, \ \int_0^T y(s) ds \le A_3, \ \int_0^T h(s) ds \le A_4$

where A_1, A_2, A_3, A_4 are positive constants independent of t and T. Then for any r > 0, $t \ge 0$,

$$y(t+r) \le \left(\frac{A_3}{r} + A_2r + A_4\right) \cdot e^{A_1A_2}.$$

Moreover,

$$\lim_{t \to +\infty} y(t) = 0.$$

Proof. See, e.g., [23].

Lemma 3.2 There holds that

$$\lim_{t \to +\infty} \|u(t) - \bar{u}\|_{H^1} = 0. \tag{3.1}$$

Proof Differentiating (1.1) with respect to x, multiplying the result by u_x , then integrating it over [0,1], by Young's inequality, we can deduce

$$\frac{d}{dt} \|u_x(t)\|^2 \le \|v_{xx}(t)\|^2 + \|u_x(t)\|^2 \le \frac{1}{2} + \frac{1}{2} \|u_x(t)\|^4 + \|v_{xx}(t)\|^2$$

which along with (1.12) and Lemma 3.1 gives

$$\lim_{t \to \infty} \|u_x(t)\|^2 = 0. \tag{3.2}$$

On the other hand, by (2.2) and the imbedding theorem, we can deduce

$$||u(t) - \bar{u}|| \le C_1 ||u_x(t)||$$

which with (3.2) gives (3.1).

Lemma 3.3 There holds that

$$\lim_{t \to +\infty} ||v(t)||_{H^1} = 0. \tag{3.3}$$

Proof Equation (1.2) can be rewritten as

$$v_t = \left(-R\frac{\theta}{u} + \int_0^x f dy\right)_x + \mu \left(\frac{v_x}{u}\right)_x. \tag{3.4}$$

Denote

$$\begin{split} \hat{p} &= \hat{p}(x,t) = R \frac{\theta}{u} - \int_0^x f dy, \\ \hat{\sigma} &= \hat{\sigma}(x,t) = -R \frac{\theta}{u} + \int_0^x f dy + \mu \frac{v_x}{u}. \end{split}$$

Then

$$v_t = \left(-\hat{p} + \mu \frac{v_x}{u}\right)_x = \hat{\sigma}_x. \tag{3.5}$$

Put

$$\hat{p}^* = \hat{p}^*(x,t) = \hat{p}(x,t) - \int_0^1 \hat{p}(x,t) dx,$$
$$\hat{\sigma}^* = \hat{\sigma}^*(x,t) = \hat{\sigma}(x,t) - \int_0^1 \hat{\sigma}(x,t) dx.$$

Then

$$\int_0^1 \hat{p}^*(x,t)dx = 0, \quad \int_0^1 \hat{\sigma}^*(x,t)dx = 0, \tag{3.6}$$

$$(\hat{p}^*)_x = \hat{p}_x, \ (\hat{\sigma}^*)_x = \hat{\sigma}_x.$$
 (3.7)

Noting (3.5) and integrating by parts, we see that

$$\|\hat{p}^*\|^2 = (\hat{p}^*, \hat{p}^*) = \left(-\hat{p}_x^*, \int_0^x \hat{p}^* dy\right)$$

$$= \left(-\hat{p}_x, \int_0^x \hat{p}^* dy\right) = \left(v_t - \mu(\frac{v_x}{u})_x, \int_0^x \hat{p}^* dy\right)$$

$$= \left(v_t, \int_0^x \hat{p}^* dy\right) - \left(\mu(\frac{v_x}{u})_x, \int_0^x \hat{p}^* dy\right)$$

$$\leq \left(\int_0^1 v_t^2 dx\right)^{1/2} \left(\int_0^1 \left(\int_0^x \hat{p}^* dy\right) dx\right)^{1/2} + \mu \int_0^1 \frac{v_x}{u} \hat{p}^* dx$$

$$\leq \varepsilon \|\hat{p}^*\|^2 + C(\varepsilon) \|v_t\|^2 + C_1(\varepsilon \|\hat{p}^*\|^2 + C(\varepsilon) \|v_x\|^2)$$

$$\leq (C_1 + 1)\varepsilon \|\hat{p}^*\|^2 + C_1(\|v_t\|^2 + \|v_x\|^2)$$
(3.8)

where (\cdot, \cdot) denotes the inner product on $L^2[0, 1]$.

Taking $\varepsilon > 0$ small enough in (3.8) and integrating it over [0, t], combining the result with (1.12), we can deduce

$$\int_0^t \|\hat{p}^*\|^2(s)ds \le C_1. \tag{3.9}$$

On the other hand, noting that $\frac{df}{dt} = f_{\xi}v + f_{t}$, we arrive at

$$\frac{d}{dt} \|\hat{p}^*\|^2 = 2(\hat{p}^*, \hat{p}_t^*) = 2(\hat{p}_x^*, -\int_0^x \hat{p}_t^* dy)
= 2(v_t, \int_0^x \hat{p}_t^* dy) - 2(\mu(\frac{v_x}{u})_x, \int_0^x \hat{p}_t^* dy)
= 2(v_t, \int_0^x \hat{p}_t^* dy) + 2\mu(\frac{v_x}{u}, \hat{p}_t^*)
\leq C_1(\|v_t\|^2 + \|v_x\|^2 + \|\hat{p}_t^*\|^2)
\leq C_1(\|v_t\|^2 + \|v_x\|^2 + \|\theta_t\|^2 + \|\frac{df}{dt}\|^2)
\leq C_1(\|v_t\|^2 \|v_x\|^2 + \|\theta_t\|^2 + \|f_{\varepsilon}\|^2 + \|f_t\|^2)$$

which combined with (3.9), Lemma 3.1 and (1.12), we can get

$$\lim_{t \to +\infty} \|\hat{p}^*(t)\|^2 = 0. \tag{3.10}$$

By (3.9) and (1.12), we have

$$\int_0^t \|\hat{\sigma}^*\|^2(s)ds \le C_1 \int_0^t (\|\hat{p}^*\|^2 + \|u_x\|^2)(s)ds \le C_1.$$
(3.11)

Noting (3.5) and integrating by parts, we see that

$$\frac{d}{dt} \|\hat{\sigma}^*\|^2 = 2(\hat{\sigma}^*, \hat{\sigma}_t^*) = 2(\hat{\sigma}_x^*, -\int_0^x \hat{\sigma}_t^* dy)
= 2(v_t, -\int_0^x \hat{\sigma}_t^* dy)
\leq C_1(\|v_t\|^2 + \|\int_0^x \hat{\sigma}_t^* dy\|^2)$$
(3.12)

where

$$\int_{0}^{x} \hat{\sigma}_{t}^{*} dy = \int_{0}^{x} \left(\hat{\sigma}_{t} - \int_{0}^{1} \hat{\sigma}_{t}(\xi, t) d\xi \right) dy$$
 (3.13)

and

$$\hat{\sigma}_{t}(x,t) = \left(-R\frac{\theta}{u} + \int_{0}^{x} f dy + \mu \frac{v_{x}}{u}\right)_{t}
= -R\frac{\theta_{t}}{u} + R\frac{\theta u_{t}}{u^{2}} + \int_{0}^{x} f_{t} dy + \mu \frac{v_{xt}}{u} - \mu \frac{v_{x}^{2}}{u^{2}}
= -R\frac{\theta_{t}}{u} + R\frac{\theta u_{t}}{u^{2}} + \int_{0}^{x} f_{t} dy + \mu \left(\frac{v_{t}}{u}\right)_{x} + \mu \frac{v_{t} u_{x} - v_{x}^{2}}{u^{2}}.$$

Noting

$$v_t(0,t) = v_t(1,t) = 0, \ \forall t > 0,$$

we deduce that for any $x \in [0, 1]$,

$$\left| \int_0^x \hat{\sigma}_t dy \right| \le C_1(\|\theta_t\| + \|v_x\| + \|f_t\|_{L^1[0,1]} + |v_t| + \|v_t\|), \tag{3.14}$$

$$\left| \int_{x}^{1} \hat{\sigma}_{t} dy \right| \leq C_{1}(\|\theta_{t}\| + \|v_{x}\| + \|f_{t}\|_{L^{1}[0,1]} + |v_{t}| + \|v_{t}\|). \tag{3.15}$$

Thus it follows from Fubini's theorem and (3.13)-(3.15) that

$$\int_{0}^{x} \hat{\sigma}_{t}^{*} dy = \int_{0}^{x} \left(\hat{\sigma}_{t} - \int_{0}^{1} \hat{\sigma}_{t}(\xi, t) d\xi \right) dy$$

$$= \int_{0}^{x} \hat{\sigma}_{t} dy - \int_{0}^{1} \int_{x}^{1} \hat{\sigma}_{t} dy d\xi$$

$$\leq C_{1}(\|\theta_{t}\| + \|v_{x}\| + \|f_{t}\|_{L^{1}[0,1]} + |v_{t}| + \|v_{t}\|)$$

which implies

$$\left\| \int_0^x \hat{\sigma}_t^* dy \right\|^2 \le C_1(\|\theta_t\|^2 + \|v_x\|^2 + \|f_t\|_{L^1[0,1]}^2 + \|v_t\|^2). \tag{3.16}$$

By (3.12) and (3.16), we deduce

$$\frac{d}{dt} \|\hat{\sigma}^*\|^2 \le C_1(\|\theta_t\|^2 + \|v_x\|^2 + \|f_t\|^2 + \|v_t\|^2)$$

which with (3.11), Lemma 3.1 and (1.12) yields

$$\lim_{t \to +\infty} \|\hat{\sigma}^*(t)\|^2 = 0. \tag{3.17}$$

Noting that

$$\int_0^1 \left(\frac{v}{u}\right)_x dx = 0,$$

we have

$$\frac{v_x}{u} = \left(\frac{v_x}{u}\right)^* + \int_0^1 \frac{v_x}{u} dx = \frac{1}{\mu} (\hat{\sigma}^* + \hat{p}^*) + \int_0^1 \frac{v u_x}{u^2} dx$$

which with Lemma 2.4 gives

$$||v_x|| \le C_1 ||\frac{v_x}{u}|| \le C_1 (||\hat{\sigma}^*|| + ||\hat{p}^*|| + ||v|| ||u_x||).$$
(3.18)

By (3.18), (3.17), (3.10) and Lemma 3.2, we derive

$$\lim_{t \to +\infty} ||v_x(t)|| = 0.$$
 (3.19)

By (1.6) and the Poincaré inequality, we deduce

$$||v(t)||_{H^1} \le C_1 ||v_x(t)||$$

which, together with (3.19), gives (3.3).

Lemma 3.4 The following estimates hold,

$$\lim_{t \to +\infty} \|\theta(t) - \bar{\theta}\|_{H^1} = 0, \quad \lim_{t \to +\infty} \|\theta(t) - \bar{\theta}\|_{L^{\infty}} = 0. \tag{3.20}$$

Proof Multiplying (1.3) by θ_t , and then integrating the resultant over [0, 1], we get

$$\frac{d}{dt} \int_{0}^{1} k \frac{\theta_{x}^{2}}{u} dx + \int_{0}^{1} c_{v} \theta_{t}^{2} dx
= \int_{0}^{1} \left(-k \frac{\theta_{x}^{2} v_{x}}{u^{2}} + (-R \frac{\theta}{u} + \mu \frac{v_{x}}{u}) v_{x} \theta_{t} + g \theta_{t} \right) dx
\leq \int_{0}^{1} (\theta_{x}^{2} |v_{x}| + \theta |v_{x} \theta_{t}| + v_{x}^{2} |\theta_{t}| + g |\theta_{t}|) dx
\leq \varepsilon \|\theta_{t}\|^{2} + C_{1}(\|v_{x}\|^{2} + \|\theta_{x}\|_{L^{4}}^{4} + \|\theta\|_{L^{\infty}[0,1]}^{2} \|v_{x}\|^{2} + \|v_{x}\|_{L^{4}}^{4} + \|g\|^{2})
\leq \varepsilon \|\theta_{t}\|^{2} + C_{1}(\|\theta_{x}\|^{3} \|\theta_{xx}\| + \|\theta_{x}\|^{4} + \|\theta\|_{L^{\infty}[0,1]}^{2} \|v_{x}\|^{2} + \|v_{x}\|^{3} \|v_{xx}\| + \|g\|^{2})
\leq \varepsilon \|\theta_{t}\|^{2} + C_{1}(\|\theta_{x}\|^{2} + \|\theta_{xx}\|^{2} + \|v_{x}\|^{2} + \|v_{xx}\|^{2} + \|g\|^{2})$$

which, by taking $\varepsilon > 0$ small enough, along with (1.12) and Lemma 3.1, we derive

$$\lim_{t \to +\infty} \|\theta_x(t)\| = 0. \tag{3.21}$$

By the Poincaré inequality, we deduce

$$\|\theta(t) - \bar{\theta}\|_{H^1} \le C_1 \|\theta_x(t)\|, \|\theta(t) - \bar{\theta}\|_{L^{\infty}[0,1]} \le C_1 \|\theta_x(t)\|$$

which, together with (3.21), gives (3.20).

4 Global Existence in H^2

In this section we will complete the proof of (1.14) and assume that the assumptions in Theorem 1.2 are valid. We begin with the following lemma. The first lemma concerns the uniform global (in time) positive lower bound (independent of t) of the absolute temperature θ .

Lemma 4.1 If $(u_0, v_0, \theta_0) \in H^1$, then the generalized global solution $(u(t), v(t), \theta(t))$ to the problem (1.1)-(1.6) satisfies

$$0 < C_1^{-1} \le \theta(x, t), \quad \forall (x, t) \in [0, 1] \times [0, +\infty). \tag{4.1}$$

Proof We use the idea in [20] to prove (4.1). If (4.1) is not true, that is, $\inf_{(x,t)\in[0,1]\times[0,+\infty)} \theta(x,t) = 0$, then there exists a sequence $(x_n,t_n)\in[0,1]\times[0,+\infty)$ such that

$$\theta(x_n, t_n) \to 0, \text{ as } n \to +\infty.$$
 (4.2)

If the sequence t_n has a subsequence, denoted also by t_n , converging to $+\infty$, then by the asymptotic behavior results in Lemma 3.4, we know that

$$\theta(x_n, t_n) \to \bar{\theta} > 0$$
, as $n \to +\infty$

which contradicts with (4.2).

If the sequence t_n is bounded, i.e., there exists a constant M > 0, independent of n, such that for any $n = 1, 2, 3, \dots, 0 < t_n \leq M$, then there exists a point $(x^*, t^*) \in [0, 1] \times [0, M]$ such that $(x_n, t_n) \to (x^*, t^*)$ as $n \to +\infty$. On the other hand, by (4.2) and the continuity of solutions, we conclude that $\theta(x_n, t_n) \to \theta(x^*, t^*) = 0$ as $n \to +\infty$, which contradicts with (2.1). Thus the proof is complete.

Lemma 4.2 There holds that

$$||v_t(t)||^2 + \int_0^t ||v_{tx}(s)||^2 ds \le C_2, \quad \forall t > 0,$$
 (4.3)

$$||v_{xx}(t)|| \le C_2, \quad \forall t > 0.$$
 (4.4)

Proof Differentiating (1.2) with respect to t, multiplying the result by v_t , then integrating it over [0, 1] and integrating by part, by (1.1), (1.7) and Young's inequality, we can derive

$$\frac{1}{2} \frac{d}{dt} \|v_{t}(t)\|^{2} + \int_{0}^{1} \mu \frac{v_{tx}^{2}}{u} dx$$

$$= \int_{0}^{1} v_{tx} (R \frac{\theta_{t}}{u} - R \frac{\theta u_{t}}{u^{2}} + \mu \frac{v_{x}^{2}}{u^{2}} - \int_{0}^{x} \frac{df}{dt} dy) dx$$

$$\leq \varepsilon \|v_{tx}\|^{2} + C_{2} (\|\theta_{t}\|^{2} + \|v_{x}\|^{2} + \|v_{x}\|_{L^{\infty}}^{2} \|v_{x}\|^{2} + \|\frac{df}{dt}\|^{2})$$

$$\leq \varepsilon \|v_{tx}\|^{2} + C_{2} (\|\theta_{t}\|^{2} + \|v_{x}\|^{2} + \|v_{tx}\|^{2} + \|f_{t}\|^{2} + \|f_{\xi}\|^{2}). \tag{4.5}$$

Integrating (4.5) over [0, t], by (1.12) and (1.9), we can derive (4.3).

Equation (1.2) can be written as

$$v_{xx} = \frac{u}{\mu} (v_t + R \frac{\theta_x}{u} - R \frac{\theta u_x}{u^2} + \mu \frac{v_x u_x}{u^2} - f). \tag{4.6}$$

By (4.6), (4.3) (1.12) and (1.9), we can get

$$||v_{xx}(t)|| = \frac{1}{\mu} \left(\int_{0}^{1} u^{2} (v_{t} + R \frac{\theta_{x}}{u} - R \frac{\theta u_{x}}{u^{2}} + \mu \frac{v_{x} u_{x}}{u^{2}} - f)^{2} dx \right)^{1/2}$$

$$\leq C_{2}(||v_{t}|| + ||\theta_{x}|| + ||\theta u_{x}|| + ||v_{x}||_{L^{\infty}} ||u_{x}|| + ||f||)$$

$$\leq \varepsilon ||v_{xx}(t)|| + C_{2}(||v_{t}(t)|| + ||v_{x}(t)|| + ||\theta_{x}(t)|| + ||u_{x}(t)|| + ||f(t)||)$$

$$\leq \varepsilon ||v_{xx}(t)|| + C_{2}$$

$$(4.7)$$

which gives (4.4).

Lemma 4.3 The following estimates hold

$$||u_{xx}(t)||^2 + \int_0^t ||u_{xx}(s)||^2 ds \le C_2, \quad \forall t > 0,$$
 (4.8)

$$\int_0^t \|v_{xxx}(s)\|^2 ds \le C_2, \quad \forall t > 0.$$
(4.9)

Proof Differentiating (1.2) with respect to x, using (1.1) $(u_{txx} = v_{xxx})$, we can deduce

$$\begin{aligned} v_{tx} &= -p_{xx} + \mu \left(\frac{v_x}{u}\right)_{xx} + f_x \\ &= -(p_u u_x + p_\theta \theta_x)_x + \mu \left(\frac{u_{xx}}{u}\right)_t - 2\mu \frac{v_{xx} u_x}{u^2} + 2\mu \frac{v_x u_x^2}{u^3} + f_x \\ &= -p_{uu} u_x^2 - p_u u_{xx} - 2p_{u\theta} \theta_x u_x - p_{\theta\theta} \theta_x^2 - p_\theta \theta_{xx} \\ &+ \mu \left(\frac{u_{xx}}{u}\right)_t - 2\mu \frac{v_{xx} u_x}{u^2} + 2\mu \frac{v_x u_x^2}{u^3} + f_x, \end{aligned}$$

i.e.

$$\mu \left(\frac{u_{xx}}{u}\right)_{t} - p_{u}u_{xx} = v_{tx} + p_{uu}u_{x}^{2} + 2p_{u\theta}\theta_{x}u_{x} + p_{\theta\theta}\theta_{x}^{2} + p_{\theta}\theta_{xx}$$

$$+2\mu \frac{v_{xx}u_{x}}{u^{2}} - 2\mu \frac{v_{x}u_{x}^{2}}{u^{3}} - f_{\xi}u. \tag{4.10}$$

Multiplying (4.10) by u_{xx}/u , and by Young's inequality Lemma 4.1, we can deduce that

$$\frac{d}{dt} \left\| \frac{u_{xx}}{u} \right\|^{2} + C_{1} \int_{0}^{1} \frac{u_{xx}^{2}}{u^{2}} dx \leq \frac{d}{dt} \left\| \frac{u_{xx}}{u} \right\|^{2} + \int_{0}^{1} R \frac{\theta}{u} \frac{u_{xx}^{2}}{u^{2}} dx$$

$$\leq \varepsilon \left\| \frac{u_{xx}}{u} \right\|^{2} + C_{1} (\|v_{tx}\|^{2} + \int_{0}^{1} (\theta^{2} u_{x}^{4} + u_{x}^{2} \theta_{x}^{2} + \theta_{xx}^{2} + \theta_{xx}^{2} + v_{xx}^{2} u_{x}^{4} + f_{\xi}^{2}) dx)$$

$$\leq \varepsilon \left\| \frac{u_{xx}}{u} \right\|^{2} + C_{1} (\|v_{tx}\|^{2} + \|u_{x}\|_{L^{\infty}[0,1]}^{2} (\|u_{x}\|^{2} + \|\theta_{x}\|^{2}) + \|\theta_{xx}\|^{2} + \|u_{x}\|_{L^{\infty}[0,1]}^{2} \|u_{x}\|_{L^{\infty}[0,1]}^{2} \|u_{x}\|^{2} + \|f_{\xi}\|^{2})$$

$$\leq \varepsilon \left\| \frac{u_{xx}}{u} \right\|^{2} + C_{2} (\|v_{tx}\|^{2} + \|u_{x}\|^{2} + \|u_{x}\| \|u_{xx}\| + \|\theta_{xx}\|^{2} + \|f_{\xi}\|^{2})$$

$$\leq 2\varepsilon \left\| \frac{u_{xx}}{u} \right\|^{2} + C_{2} (\|v_{tx}\|^{2} + \|u_{x}\|^{2} + \|\theta_{xx}\|^{2} + \|f_{\xi}\|^{2})$$

$$(4.11)$$

which, by taking $\varepsilon > 0$ small enough and using Lemmas 2.1-2.8, Lemmas 4.1-4.2, gives (4.8).

Differentiating (1.2) with respect to x, by (1.1), we can derive

$$v_{xxx} = \frac{u}{\mu} \left(v_{tx} + R \frac{\theta_{xx}}{u} - 2R \frac{\theta_{x}u_{x}}{u^{2}} - R \frac{\theta u_{xx}}{u^{2}} + 2R \frac{\theta u_{x}^{2}}{u^{3}} + \mu \frac{u_{xx}u_{t}}{u^{2}} + 2\mu \frac{v_{xx}u_{x}}{u^{2}} - 2\mu \frac{v_{x}u_{x}^{2}}{u^{3}} + f_{x} \right).$$

Thus

$$\int_{0}^{t} \|v_{xxx}(s)\|^{2} ds \leq C_{2} \int_{0}^{t} \int_{0}^{1} (v_{tx}^{2} + \theta_{xx}^{2} + \theta_{x}^{2} u_{x}^{2} + \theta^{2} u_{xx}^{2} + \theta^{2} u_{x}^{4} + \theta^{2} u_{xx}^{4} +$$

$$\leq C_{2} \int_{0}^{t} (\|v_{tx}\|^{2} + \|u_{xx}\|^{2} + \|u_{x}\|^{2} + \|f_{\xi}\|^{2}) ds$$

$$+ C_{2} \int_{0}^{t} (\|u_{x}\|_{L^{\infty}}^{2} (\|\theta_{x}\|^{2} \|v_{xx}\|^{2})$$

$$+ \|v_{x}\|_{L^{\infty}}^{2} \|u_{xx}\|^{2} + \|u_{x}\|_{L^{\infty}}^{2} \|v_{x}\|_{L^{\infty}}^{2} \|u_{x}\|^{2}) ds$$

$$\leq C_{2} \int_{0}^{t} (\|v_{tx}\|^{2} + \|u_{xx}\|^{2} + \|u_{x}\|^{2} + \|f_{\xi}\|^{2}) ds$$

$$+ C_{2} \int_{0}^{t} (\|v_{x}\|_{L^{\infty}}^{2} + \|u_{x}\|_{L^{\infty}}^{2}) ds$$

$$\leq C_{2}$$

which gives (4.9).

Lemma 4.4 The following estimates hold,

$$\|\theta_t(t)\|^2 + \int_0^t \|\theta_{xt}\|^2 ds \le C_2, \quad \forall t > 0,$$
 (4.12)

$$\|\theta_{xx}(t)\| \le C_2, \quad \forall t > 0,$$
 (4.13)

$$\int_{0}^{t} \|\theta_{xxx}(s)\|^{2} \le C_{2}, \quad \forall t > 0.$$
(4.14)

Proof Differentiating (1.3) with respect to t and multiplying the resultant by θ_t , we can derive

$$\frac{c_{v}}{2} \frac{d}{dt} \|\theta_{t}(t)\|^{2} + \int_{0}^{1} k \frac{\theta_{xt}^{2}}{u} dx \\
= \int_{0}^{1} \left(k \frac{\theta_{x} \theta_{xt} v_{x}}{u^{2}} - R \frac{\theta_{t}^{2} v_{x} + \theta \theta_{t} v_{xt}}{u^{2}} + R \frac{\theta v_{x}^{2} \theta_{t}}{u^{2}} \right) \\
+ \mu \frac{2v_{x} v_{xt} \theta_{t}}{u} - \mu \frac{v_{x}^{3} \theta_{t}}{u^{2}} + \frac{dg}{dt} \theta_{t} dx \\
\leq \varepsilon \|\theta_{xt}\|^{2} + C_{2} \int_{0}^{1} (\theta_{x}^{2} v_{x}^{2} + |\theta_{t}^{2} v_{x}| + |\theta \theta_{t} (v_{xt} + v_{x}^{2})| + |v_{x} v_{xt} \theta_{t}| + |v_{x}^{3} \theta_{t}| + |\frac{dg}{dt} \theta_{t}|) dx \\
\leq \varepsilon \|\theta_{xt}\|^{2} + C_{2} \left(\|v_{x}\|_{L^{\infty}}^{2} \|\theta_{x}\|^{2} + \|\theta_{t}\|_{L^{4}}^{4} + \|v_{x}\|^{2} + \|\theta\|_{L^{\infty}} (\|v_{tx}\|^{2} + \|v_{x}\|_{L^{4}}^{4} + \|\theta_{t}\|^{2}) \\
+ \|v_{x}\|_{L^{\infty}}^{2} (\|v_{xt}\|^{2} + \|\theta_{t}\|^{2}) + \|v_{x}\|_{L^{\infty}} (\|v_{x}\|_{L^{4}}^{4} + \|\theta_{t}\|^{2}) + \|g_{t}\|^{2} + \|g_{\xi}\|^{2} + \|\theta_{t}\|^{2}) \\
\leq \varepsilon \|\theta_{xt}\|^{2} + C_{2} (\|\theta_{x}\|^{2} + \|v_{x}\|^{2} + \|v_{tx}\|^{2} + \|\theta_{t}\|^{2}) + C_{2} (\|v_{x}\|_{L^{4}}^{4} + \|\theta_{t}\|_{L^{4}}^{4}) \\
+ C_{2} (\|g_{t}\|^{2} + \|g_{\xi}\|^{2}) \\
\leq \varepsilon \|\theta_{xt}\|^{2} + C_{2} (\|\theta_{x}\|^{2} + \|v_{x}\|^{2} + \|v_{tx}\|^{2} + \|\theta_{t}\|^{2}) + C_{2} (\|g_{t}\|^{2} + \|g_{\xi}\|^{2}) \\
\leq \varepsilon \|\theta_{xt}(t)\|^{2} + C_{2} (\|\theta_{x}\|^{2} + \|v_{x}\|^{2} + \|v_{tx}\|^{2} + \|\theta_{t}\|^{2}) + C_{2} (\|g_{t}\|^{2} + \|g_{\xi}\|^{2}). \tag{4.15}$$

Taking $\varepsilon > 0$ small enough, integrating (4.15) over [0, t], and using Lemmas 2.1-2.8, Lemmas 4.1-4.3, we can get (4.12).

(1.3) can be rewritten as

$$\theta_{xx} = \frac{u}{k} \left(c_v \theta_t + k \frac{\theta_x u_x}{u^2} + R \frac{\theta v_x}{u} - \mu \frac{v_x^2}{u} - g \right).$$

Thus

$$\|\theta_{xx}(t)\|^{2} \leq C_{2}(\|\theta_{t}\|^{2} + \|\theta_{x}u_{x}\|^{2} + \|\theta v_{x}\|^{2} + \|v_{x}\|_{L^{4}[0,1]}^{4} + \|g\|^{2})$$

$$\leq C_{2}(\|\theta_{t}\|^{2} + \|\theta_{x}\|_{L^{\infty}[0,1]}^{2} \|u_{x}\|^{2} + \|\theta\|_{L^{\infty}[0,1]}^{2} \|v_{x}\|^{2} + \|v_{x}\|^{2} + \|v_{x}\|^{3} \|v_{xx}\| + \|g\|^{2})$$

$$\leq \varepsilon \|\theta_{xx}\|^{2} + C_{2}(\|\theta_{t}\|^{2} + \|\theta_{x}\|^{2} + \|v_{x}\|^{2} + \|g\|^{2})$$

$$(4.16)$$

which, combine with (4.12), implies

$$\|\theta_{xx}(t)\| \le C_2(\|\theta_t(t)\| + 1) \le C_2.$$

Differentiating (1.3) with respect to x, we can derive

$$\theta_{xxx} = \frac{u}{k} \left(c_v \theta_{tx} + k \frac{\theta_{xx} u_x}{u^2} + (k \frac{\theta_x u_x}{u^2} + R \frac{\theta v_x}{u} - \mu \frac{v_x^2}{u} - g)_x \right)$$

$$= \frac{u}{k} \left(c_v \theta_{tx} + k \frac{\theta_{xx} u_x}{u^2} + k \frac{\theta_{xx} u_x + \theta_x u_{xx}}{u^2} \right)$$

$$-2k \frac{\theta_x u_x^2}{u^3} - R \frac{\theta_x v_x + \theta v_{xx}}{u} + R \frac{\theta v_x u_x}{u^2} + 2\mu \frac{v_x v_{xx}}{u} - \mu \frac{v_x^2 u_x}{u^2} + g_x \right). \quad (4.17)$$

Integrating (4.17) over $[0,1] \times [0,t]$, we get

$$\int_{0}^{t} \|\theta_{xxx}(s)\|^{2} ds \leq C_{2} \int_{0}^{t} (\|\theta_{tx}\|^{2} + \|u_{x}\|_{L^{\infty}}^{2} \|\theta_{xx}\|^{2} + \|\theta_{x}\|_{L^{\infty}}^{2} \|u_{xx}\|^{2}
+ \|\theta_{x}\|_{L^{\infty}}^{2} \|u_{x}\|_{L^{\infty}}^{2} \|u_{x}\| + \|v_{x}\|_{L^{\infty}}^{2} \|\theta_{x}\|^{2} + \|\theta\|_{L^{\infty}}^{2} \|v_{xx}\|^{2}
+ \|\theta\|_{L^{\infty}}^{2} \|v_{x}\|_{L^{\infty}}^{2} \|u_{x}\|^{2} + \|v_{x}\|_{L^{\infty}}^{2} \|v_{xx}\|^{2}
+ \|v_{x}\|_{L^{\infty}}^{2} \|u_{x}\|_{L^{\infty}}^{2} \|v_{x}\|^{2} + \|g_{\xi}\|^{2}) ds$$

$$\leq \int_{0}^{t} (\|\theta_{tx}\|^{2} + \|\theta_{xx}\|^{2} + \|u_{xx}\|^{2}
+ \|u_{x}\|^{2} + \|\theta_{x}\|^{2} + \|v_{xx}\|^{2} + \|v_{x}\|^{2} + \|g_{\xi}\|^{2}) ds$$

which, along with (4.8), (4.12), Lemmas 2.1-2.8, Lemmas 4.1-4.3, gives estimate (4.14).

5 Asymptotic Behavior in H^2

In this section we will complete the proof of (1.15) and assume that the assumptions in Theorem 1.2 are valid.

Lemma 5.1 The following estimate holds

$$\lim_{t \to +\infty} \|u(t) - \bar{u}\|_{H^2} = 0 \tag{5.1}$$

where $\bar{u} = \int_0^1 u(x, t) dx = \int_0^1 u_0(x) ds$.

Proof Differentiating (1.1) with respect to twice x, multiplying the result by u_{xx} , then integrating it over [0, 1], by Young's inequality, we can deduce

$$\frac{d}{dt}||u_{xx}||^2 \le ||v_{xxx}||^2 + ||u_{xx}||^2 \le \frac{1}{2} + \frac{1}{2}||u_{xx}||^4 + ||v_{xxx}||^2$$

which, along with Lemmas 3.1 and 4.3, gives

$$\lim_{t \to +\infty} ||u_{xx}(t)||^2 = 0. \tag{5.2}$$

Lemma 5.2 The following estimate holds

$$\lim_{t \to +\infty} ||v(t)||_{H^2} = 0. \tag{5.3}$$

Proof Noting

$$\frac{d}{dt}\|f\|^2 = \frac{d}{dt} \int_0^1 f^2 dx = 2 \int_0^1 f \cdot \frac{df}{dt} dx \le C_2(\|f\|^2 + \|f_t\|^2 + \|f_\xi\|^2). \tag{5.4}$$

By (5.4), (1.9) and Lemma 3.1, we can conclude

$$\lim_{t \to +\infty} ||f(t)||^2 = 0. \tag{5.5}$$

Similarly, we can get

$$\lim_{t \to +\infty} \|g(t)\|^2 = 0. \tag{5.6}$$

By (3.4), (3.5), we have

$$\hat{p}_x = \left(R\frac{\theta}{u}\right)_x - f = R\frac{\theta_x}{u} - R\frac{\theta u_x}{u^2} - f$$

$$= -v_t + \mu \left(\frac{v_x}{u}\right)_x. \tag{5.7}$$

Hence

$$\|\hat{p}_x(t)\| \le C_2(\|\theta_x\| + \|u_x\| + \|f\|)$$

which implies

$$\lim_{t \to +\infty} \|\hat{p}_x(t)\|^2 = 0. \tag{5.8}$$

Differentiating (5.7) with respect to t, multiplying the result by v_t , then integrating it over [0, 1], and integrating by parts, using (1.1), (1.7)-(1.10), we can get

$$\frac{1}{2} \frac{d}{dt} \|v_{t}(t)\|^{2} + \int_{0}^{1} \mu \frac{v_{xt}^{2}}{u} dx = \int_{0}^{1} \hat{p}_{t} v_{tx} dx + \int_{0}^{1} \mu \frac{v_{x}^{2} v_{tx}}{u^{2}} dx \\
\leq \varepsilon \|v_{tx}\|^{2} + \|\hat{p}_{t}\|^{2} + \|v_{x}\|_{L^{4}}^{4} \\
\leq \varepsilon \|v_{tx}\|^{2} + \|\theta_{t}\|^{2} + \|\frac{df}{dt}\|^{2} + \|\theta\|_{L^{\infty}}^{2} \|v_{x}\|^{2} + \|v_{x}\|^{3} \|v_{xx}\| + \|v_{x}\|^{4}.$$

Thus for small $\varepsilon > 0$,

$$\frac{d}{dt}\|v_t(t)\|^2 + \|v_{tx}(t)\|^2 \le C_2(\|\theta_t(t)\|^2 + \|v_x(t)\|^2 + \|f_t\|^2 + \|f_\xi\|^2). \tag{5.9}$$

Hence we infer from (5.9), (1.14), (1.9) and Lemma 3.1,

$$\lim_{t \to +\infty} ||v_t(t)||^2 = 0. \tag{5.10}$$

By (5.7) and (1.14), we can conclude

$$||v_{xx}(t)|| \le C_2(||v_t(t)|| + ||\hat{p}_x(t)|| + ||v_x(t)||)$$

which with (5.10), (5.8), (1.13) gives

$$\lim_{t \to +\infty} ||v_{xx}(t)||^2 = 0. \tag{5.11}$$

Thus (5.3) follows from (5.11) and (1.13).

Lemma 5.3 There holds that

$$\lim_{t \to +\infty} \|\theta(t) - \bar{\theta}\|_{H^2} = 0. \tag{5.12}$$

Proof By (4.15), we can get

$$\frac{d}{dt} \|\theta_t(t)\|^2 + \|\theta_{xt}(t)\|^2
\leq C_2(\|\theta_x\|^2 + \|v_x\|^2 + \|v_{tx}\|^2 + \|\theta_t\|^2) + C_2\|\frac{dg}{dt}\|^2
\leq C_2(\|\theta_x(t)\|^2 + \|v_x(t)\|^2 + \|v_{tx}(t)\|^2 + \|\theta_t(t)\|^2) + C_2(\|g_t\|^2 + \|g_{\xi}\|^2)$$
(5.13)

which combines with (1.14) and Lemma 3.1, we can conclude

$$\lim_{t \to +\infty} \|\theta_t(t)\| = 0. \tag{5.14}$$

By (4.16), we get

$$\|\theta_{xx}(t)\|^2 \le C_2(\|\theta_t(t)\|^2 + \|\theta_x(t)\|^2 + \|v_x(t)\|^2 + \|g\|^2)$$
(5.15)

which with (5.6), (5.13), (1.13) gives

$$\lim_{t \to +\infty} \|\theta_{xx}(t)\| = 0. \tag{5.16}$$

Thus (5.12) follows from (5.16) and (1.13).

ACKNOWLEDGEMENTS This work was supported in part by grants of the NNSF of China (No. 10571024) and of the Institute of Mathematical Sciences (IMS), The Chinese University of Hong Kong. The authors also take this opportunity to thank the referee for his/her helpful suggestions which have improved the original version of this paper.

References

- [1] A. A. Amosov and A. A. Zlotnik, Global generalized solutions of the wquations of the one-dimensional motion of a viscous heat-conducting gas, Soviet Math. Dokl. Vol. 38, No. 1, (1989).
- [2] S. Jiang, on the asymptotic behavior of the motion of a viscous, heat-conducting, one-dimensional real gas, Math. Z., 216(1994), 317-336.
- [3] S. Jiang, Large-time behavior of solutions to the equations of a viscous polytropic ideal gas, Ann. Mat. Pura, Appl. CLXXV (1998), 253-275.
- [4] S. Kawashima, T.Nishida, Global solutions to the initial boundary value problems for the equations of one-dimensional motion of viscous polytropic gases, J. Math. Kyoto Univ. 21 (1981), 825-837.
- [5] A. V. Kazhikhov, Sur la solubilite globale des problemes mono dimensoinnels aux valeurs initiales-limites pour les equations dun gaz visqueux et calorifre. C. R. Acad. Sci. Paris Ser. A 284 (1977), 317-320.
- [6] A. V. Kazhikhov and V. V. Shelukhin, Unique global solution with respect to time of initial boundary value problems for one-dimensional equations of a viscous gas, J. Appl. Math. Mech. 41 (1977), 273-282.
- [7] A. V. Kazhikhov, To a theory of boundary value problems for equations of one-dimensional nonstationary motion of viscous heat-conduction gases, in Boundary Value Problems for Hydrodynamical Equations, No. 50, Institute of Hydrodynamics, Siberian Branch Acad. USSR, 1981, 37-62 (in Russian).
- [8] P. B. Mucha, Compressible Navier-Stokes system in 1-D, Math. Meth. Appl. Sci. 24 (2001), 607-622.
- [9] T. Nagasawa, On the one-dimensional motion of the polytropic ideal gas non-fixed on the boundary, J. Differential Equations 65 (1986), 49-67.
- [10] T. Nagasawa, On the outer pressure problem of the one-dimensional polytropic ideal gas, Japan J. Appl. Math. 5 (1988), 53-85.
- [11] T. Nagasawa, Global asymptotics of the outer pressure problem with free boundary, Japan J. Appl. Math. 5 (1988), 205-224.
- [12] T. Nagasawa, On the asymptotic behaviour of the one-dimensional motion of the polytropic ideal gas with stress-free condition, Quart. Appl. Math. 46 (1988), 665-679.
- [13] M. Okada and S. Kawashima, On the equations of one-dimensional motion of compressible viscous fluids, J. Math. Kyoto Univ. 23 (1983), 55-71.

- [14] Y. Qin, Global existence and asymptotic behaviour of solutions to nonlinear hyperbolic-parabolic coupled systems with arbitrary initial data, Ph.D. thesis, Fudan University (1998).
- [15] Y. Qin, Global existence and asymptotic behaviour of solutions to a system of equations for a nonlinear one-dimensional viscous, heat-conducting real gas, Chinese Ann. Math. Ser. A 20 (1999), 343-354 (in Chinese).
- [16] Y. Qin, Global existence and asymptotic behavior for the solutions to nonlinear viscous, heat-conductive, one-dimensional real gas, Adv. Math. Sci. Appl., 10 (2000), 129-148.
- [17] Y. Qin, Global existence and asymptotic behaviour for a viscous, heat-conducting one-dimensional real gas with fixed and thermally insulated end points, Nonlinear Analysis, 44 (2001), 413-441.
- [18] Y. Qin, Global existence and asymptotic behavior of solution to the system in one-dimensional nonlinear thermoviscoelasticity, Quart. Appl. Math., 59 (2001), 113-142.
- [19] Y. Qin, Global existence and asymptotic behaviour for a viscous, heat-conductive, one-dimensional real gas with fixed and constant temperature boundary conditions, Advances in Differential Equations, 7 (2002), 129-154.
- [20] Y. Qin, Exponential stability for a nonlinear one-dimensional heat conductive viscous real gas, Journal of Mathematical Analysis and Applications, 272 (2002), 507-535.
- [21] Y. Qin and J. E. M. Rivera, Universal attractor for a nonlinear one-dimensional heat conductive viscous real gas, Proceedings of the Royal Society of Edinburgh A, 132 (2002), 685-709..
- [22] Y. Qin and J. E. M. Rivera, Exponential stability and universal attractors for the Navier-Stokes equations of compressible fluids between two horizontal parallel plates in \mathbb{R}^3 , Applied Numerical Mathematics, 47 (2) (2003), 249-256.
- [23] Y. Qin and Y. Zhao, Global existence and asymptotic behavior of compressible Navier-Stokes equations for a 1-D isothermal viscous gas, Math. Mode. Meth. Appl. Sci., 18(10) (2008), to appear.
- [24] Y. Qin, Y. Wu and F. Liu, On the Cauchy problem for a one-dimensional compressible viscous polytropic ideal gas, Portugaliae Math, 64 (2007), 87-126.
- [25] A. Tani, On the first initial-boundary value problem of compressible viscous fluid motion, Publ. Res. Inst. Math. Sci. 13 (1977), 193-253.
- [26] Shigenori Yanagi, Existence of periodic solutions for a one-dimensional isentropic model system of compressible viscous gas, Nonlinear Analysis 46 (2001), 279-298.
- [27] Ting Zhang and Daoyuan Fang, Global behavior of compressible Navier-Stokes equations with denerate viscosity coefficient, Preprint.

- [28] S. Zheng and Y. Qin, Maximal attractor for the system of one-dimensional polytropic viscous ideal gas, Quart. Appl. Math. 59 (2001), 579-599.
- [29] A. A. Zlotnik and Nguyen Zha Bao, Properties and asymptotic behavior of solutions of some problems of one-dimensional of a viscous barotropic gas, Mathematical Notes, Vol.55, No.5, 1994.
- [30] A. A. Zlotnik and A. A. Amosov, On stability of generalized solutions of the equations of one-dimensional motion of a ciscous heat conducting gas, Siberian Mathematical Journal, Vol. 38, No. 4, 1997.